On Twinning in Smectic Crystals

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It is shown that mechanical twinning in smectic crystals is possible. The structure of the boundary of twins for a small disorientation of crystallites is determined. The periodic twin structure, which should appear at the tension of the smectic layer, is proposed.

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The energy of small deformations of a smectic crystal is given by the expression [1]

$$\mathscr{E} = \int \frac{A}{2} \left\{ \left(\partial_z u - \frac{\left(\partial_\alpha u \right)^2}{2} \right)^2 + \lambda^2 \left(\Delta_\perp u \right)^2 \right\} dV, \qquad (1)$$

where *u* is the displacement of the layers along the *z* axis (in the initial homogeneous undeformed state of the smectic crystal, layers lie in the *xy* plane), *A* is the elastic modulus, λ is the length parameter, ∂_{α} is the gradient vector in the *xy* plane, and $\Delta_{\perp} = \partial_{\alpha}^2$. According to Eq. (1), the undeformed state turned by small angle $\theta \ll 1$ in the *xz* plane (in this case, $\partial_x u = \theta$) corresponds to the derivative $\partial_z u = \theta^2/2$.

Let us consider the boundary between the states $\partial_x u = \pm \Theta \ (x \longrightarrow \pm \infty)$ that lies in the *yz* plane. The quantity $\partial_z u$ is unchanged inside the boundary. The variation equilibrium equation in the problem under consideration reduces to the form

$$\lambda^2 f''' + \frac{\theta^2}{2} f' - \frac{3}{2} f^2 f' = 0, \qquad (2)$$

where $f = \partial_x u$. The solution of this equation has the form $f = \theta \tanh(qx/2\lambda)$. The energy of the unit area of this boundary is given by the formula

$$\sigma = \frac{2}{3}A\lambda\theta^3.$$
 (3)

The twin structure of smectic crystals must be observed under the conditions of Helfrich instability at strains noticeably larger than the critical value (see the problem in [1, Sect. 45]: the smectic layer of thickness L bounded by solid walls parallel to the smectic layer is extended along the z axis). At very small tensions $\delta L >$

 $\delta L_c = 2\pi\lambda$, i.e., when δL is about the smectic period, the homogeneous state becomes unstable with respect to the appearance of a periodic structure in the *xy* plane with wavenumber $k_c = \sqrt{\pi/\lambda L}$. At a much larger strain $\delta L \gg \delta L_c$, a twin structure as that schematically shown in the figure should appear. The parameters of this structure are determined by minimizing the total energy of twin boundaries ($\theta = \varepsilon$ at the vertical boundaries and $\theta = \varepsilon/2$ at the boundaries of the triangular regions). According to geometric consideration, the angle at the vertex of a triangle is equal to ε , the height *H* of the triangles is related to the structure period *d* as $\tan(\varepsilon/2) =$ d/2H, and the quantity δL is related to the parameter ε as $\delta L = (L - H)(\cos^{-1}\varepsilon - 1)$. The energy density of the structure proposed above is given by the expression

$$\frac{1}{Ld} \left\{ 2 \frac{L-H}{\cos\varepsilon} \sigma(\varepsilon) + 4 \frac{H}{\cos(\varepsilon/2)} \sigma(\varepsilon/2) \right\}.$$
(4)

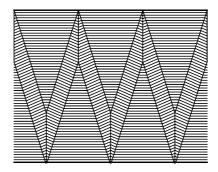


Figure.

731

In view of the indicated geometric relationships, energy (4) is a function of one parameter ε at a given tension $\delta L/L$. When $\delta L \ll L$, angle ε is small. In this case, with the use of result (3), the minimum of energy (4) is found to correspond to $\varepsilon = \sqrt{6\delta L/L}$. In this case, H = 2L/3 and $d = 2\sqrt{2L\delta L/3}$. I am grateful to E.I. Kats for stimulating discussions.

REFERENCES

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