

The Cryogenic Nanosensors Workshop, Björkliden, Kiruna, Sweden, March
20-27, 2011

The peculiar THz features of surface electromagnetic waves

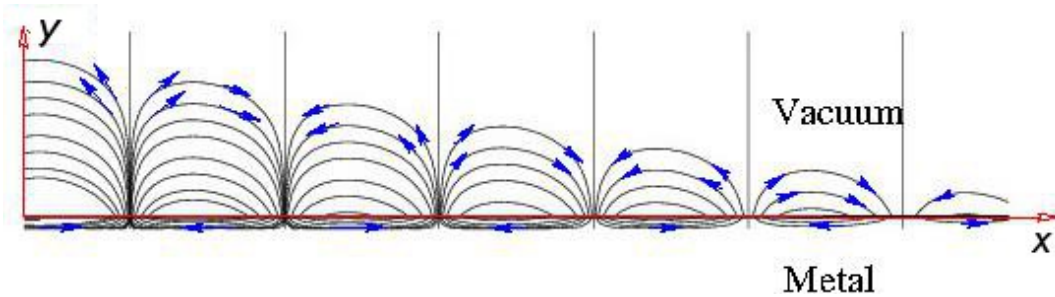
V.V. Zavyalov

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P.L. Kapitza Institute for Physical Problems Kosygina 2, Moscow, Russia;

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- V.S. Edelman, I.I. Smol'yaninov, and V.V. Zav'yalov, *Spectroscopic measurements of light emitted by scanning tunneling microscope.*
Phys. Lett. **A158**, pp.337 (1991)
 - G.D. Bogomolov, Uk Young Jeong, G.N. Zhizhin, A.K. Nikitin, V.V. Zavyalov, G.M. Kazakevich and Byung Cheol Lee,
Generation of surface electromagnetic waves in terahertz spectral range by free-electron laser radiation and their refractive index determination.
Nucl. Instr. & Meth. in Phys. Res. **A 543** p.96 (2005)
 - G.N. Zhizhin, A.K. Nikitin, G.D. Bogomolov, V.V. Zavyalov, Jeong Young Uk, Lee Byung Cheol, Seong Hee Park and Hyuk Jin Cha,
Absorption of surface plasmons in a metal-cladding layer-air structure in the terahertz frequency range.
Optics and Spectroscopy **100**, No.5, pp.734-738 (2006)
 - Zhizhin G.N., Nikitin A.K., Bogomolov G.D., Zavyalov V.V., Jeong Y.U., Lee B.C., Park S.H., Cha H.J.,
Aluminum optical constants in the far infrared determined from surface electromagnetic waves characteristics.
Proc. SPIE, 2006, v.6162, Art. 61620C.

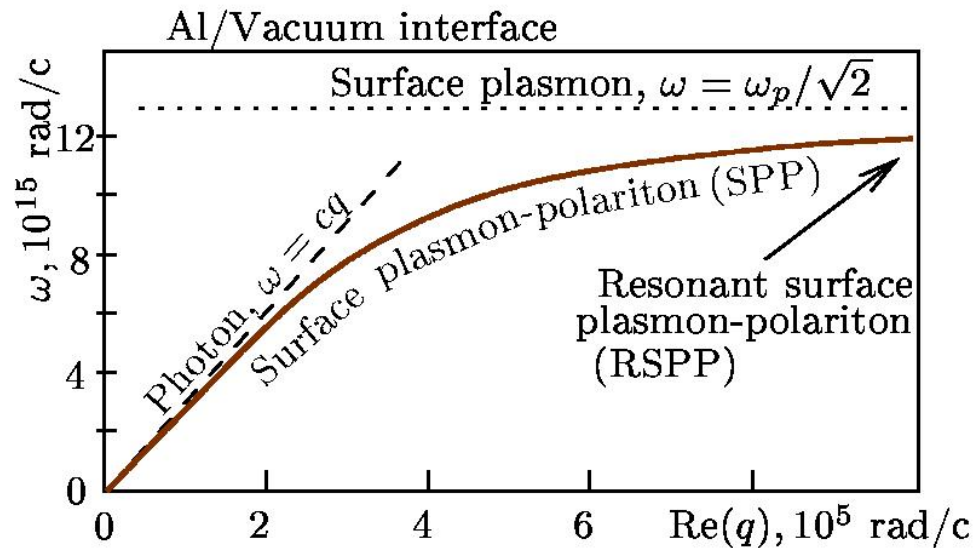
SEW is an evanescent-field solution of Maxwell's equations

$$\text{rot } \mathbf{E} = i \frac{\omega}{c} \mathbf{B}, \quad \text{rot } \mathbf{B} = \frac{4\pi}{c} \sigma \mathbf{E} - i \frac{\omega}{c} \varepsilon \mathbf{E}, \quad \text{div } \mathbf{B} = 0, \quad \varepsilon = \varepsilon' + i\varepsilon''$$



$$E_{\text{vac}} \sim e^{i\omega t + qx + k_{\text{vac}} z}$$

$$E_{\text{met}} \sim e^{i\omega t + qx + k_{\text{met}} z}$$



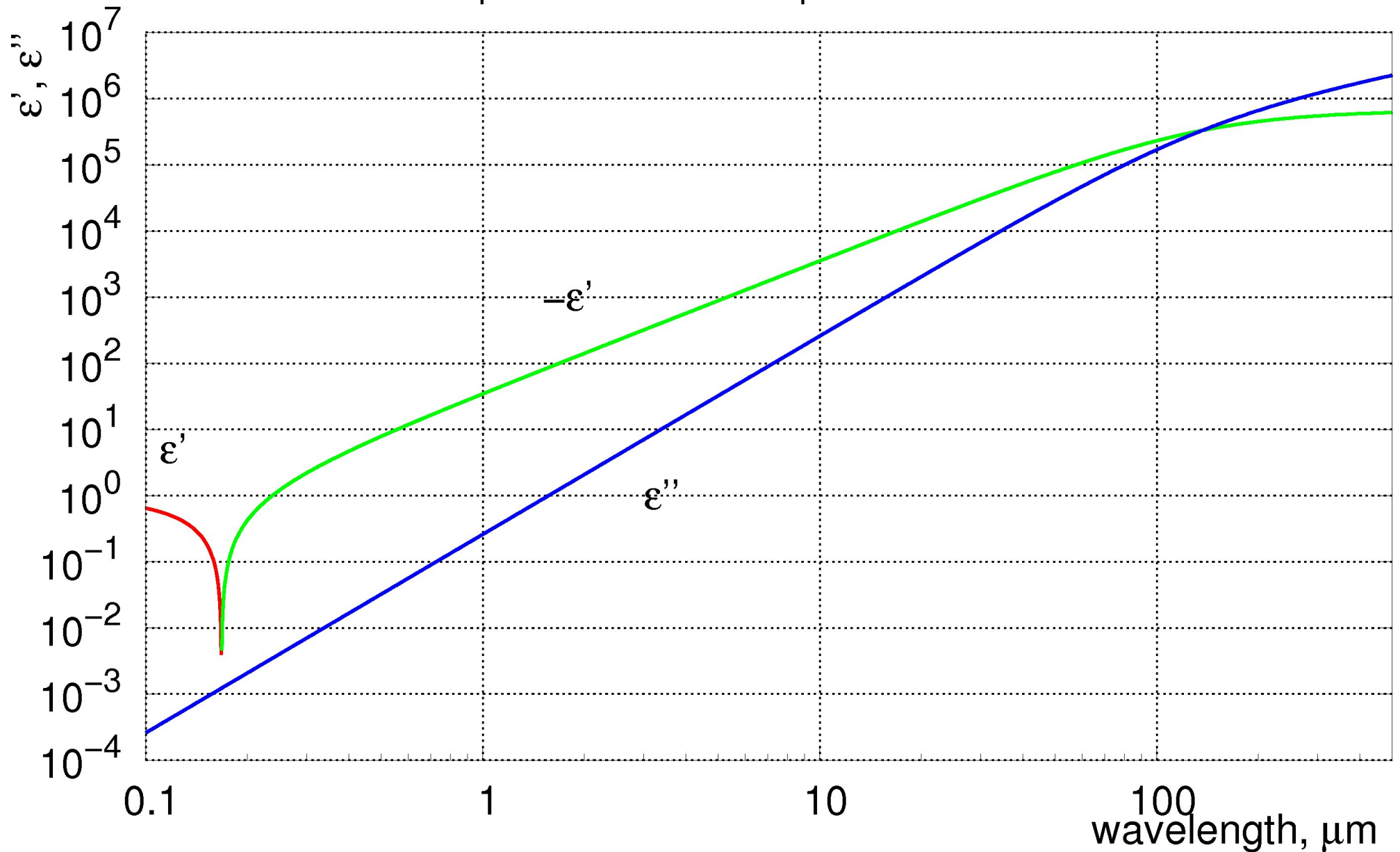
$$q = k_0 \sqrt{\frac{\varepsilon}{1 + \varepsilon}}$$

$$k_{\text{vac}} = \frac{q}{\sqrt{\varepsilon}}$$

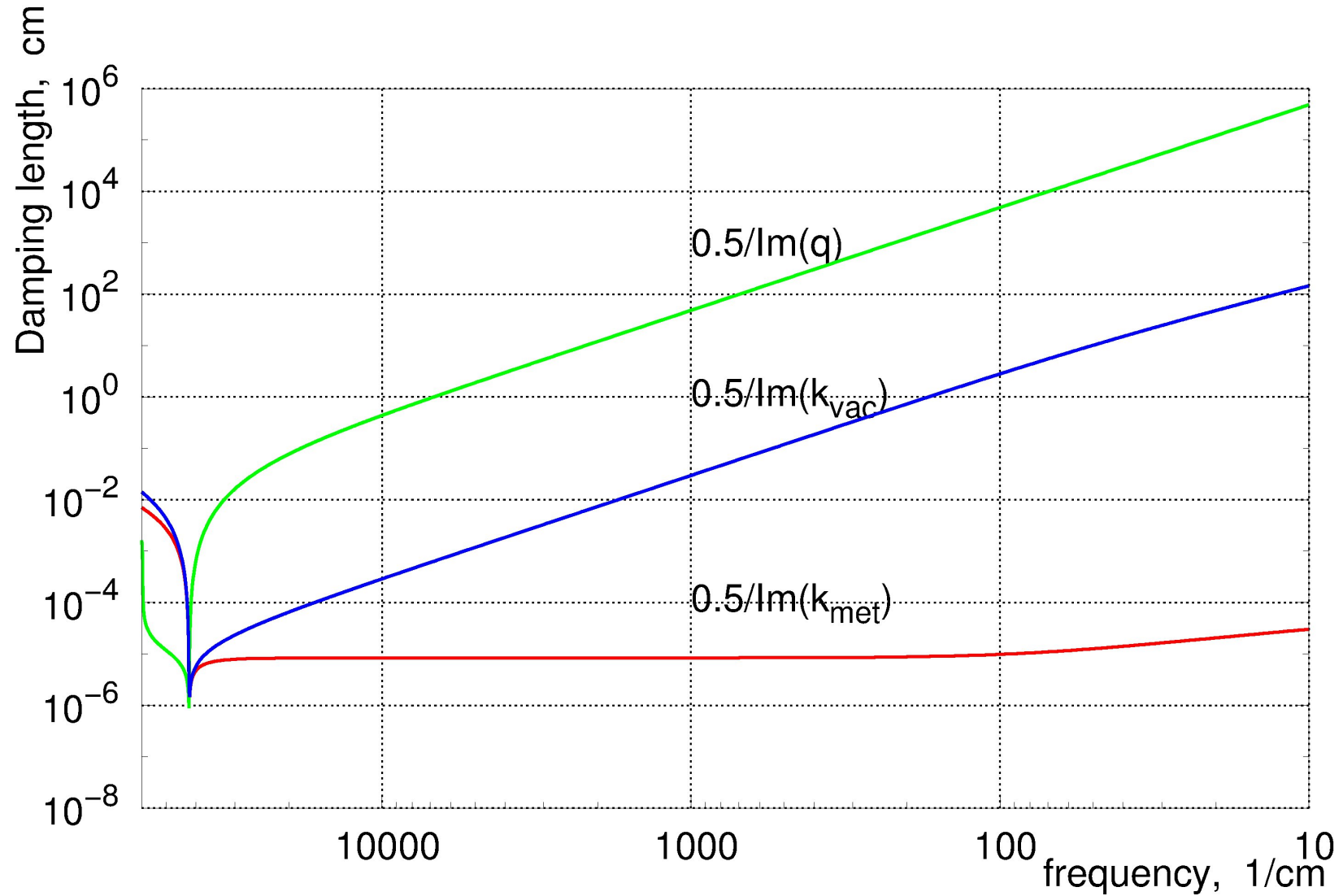
$$k_{\text{met}} = \sqrt{\varepsilon} q$$

Dielectric function for Cu (Drude model),

$$\epsilon = \epsilon' + i\epsilon'' = 1 - \frac{v_p^2}{v^2 + ivv_\tau}, \quad v_p = 59600 \text{ cm}^{-1}, \quad v_\tau = 73.2 \text{ cm}^{-1}$$

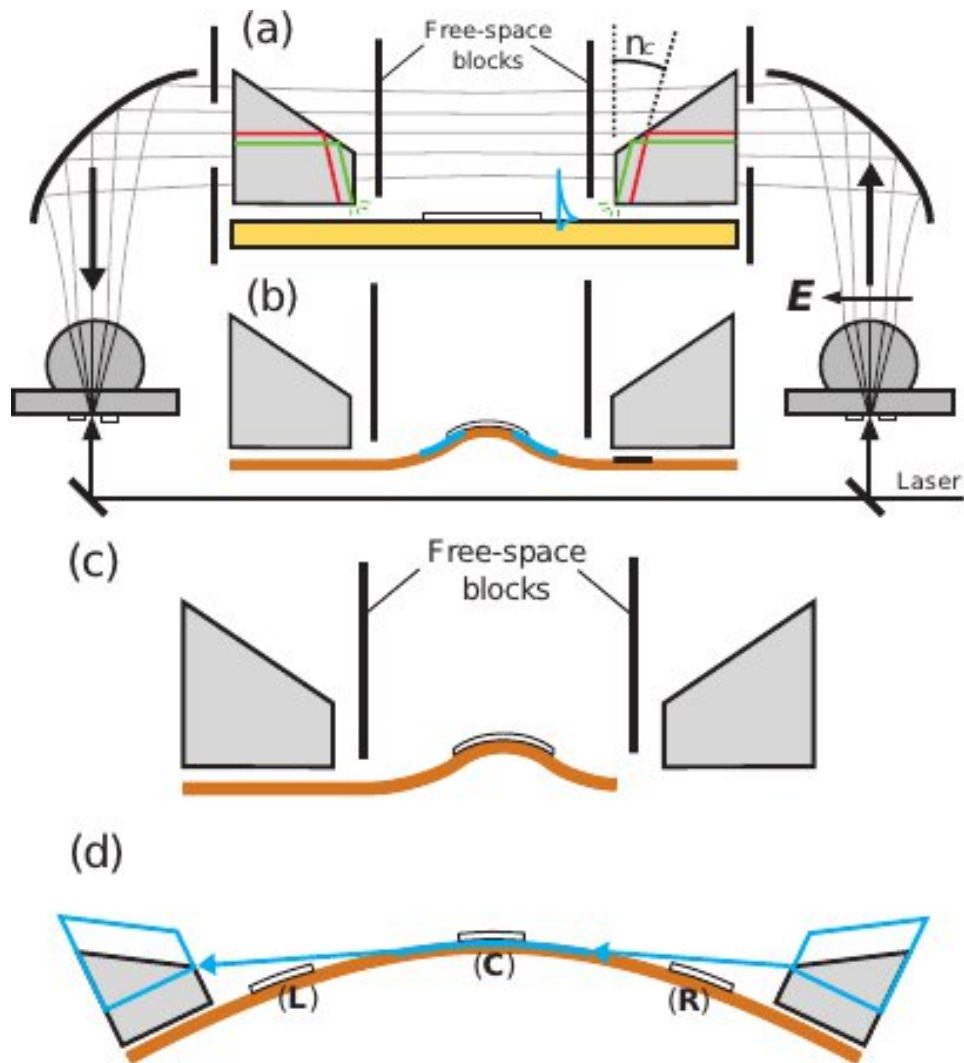


Damping lengths for Al

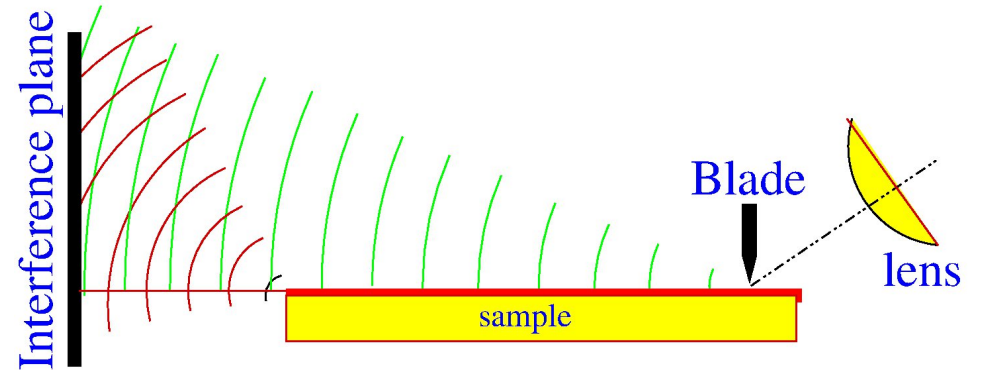


experimental observation of THz SEW

J F. O'Hara, R D. Averitt, A. J. Taylor
Opt.Express **13**, 6117 (2005)

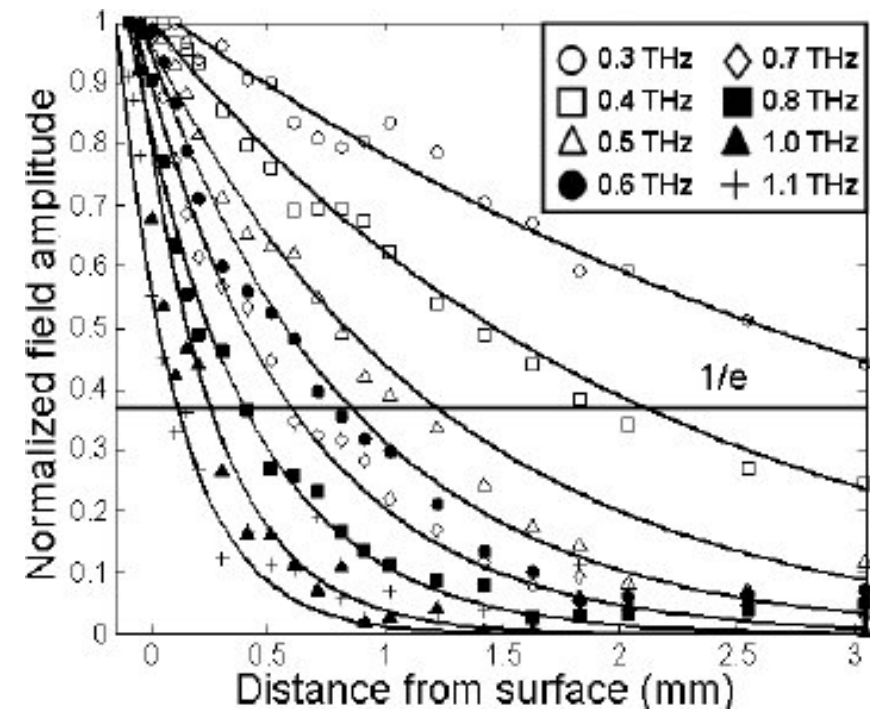
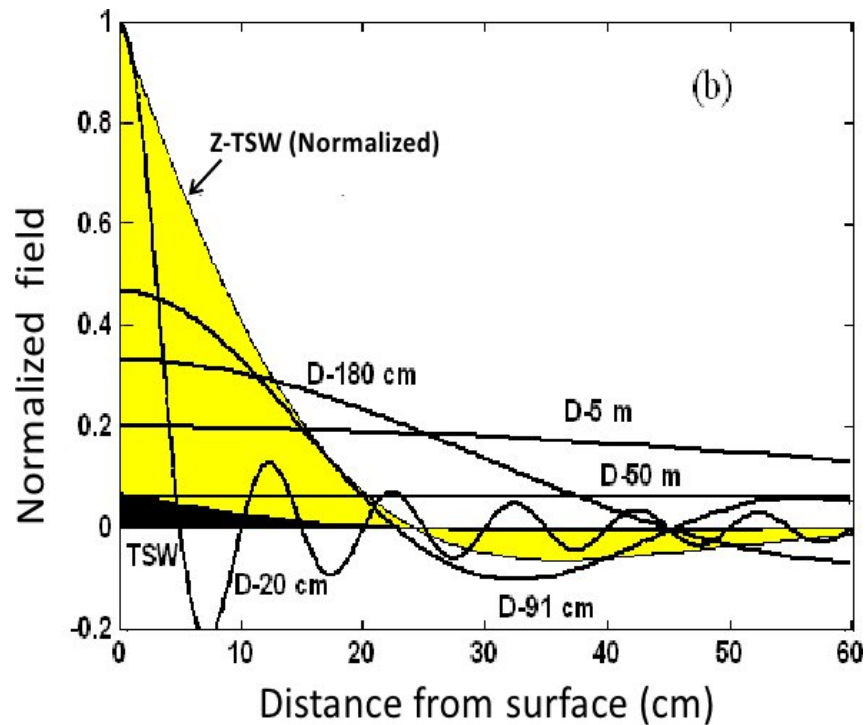
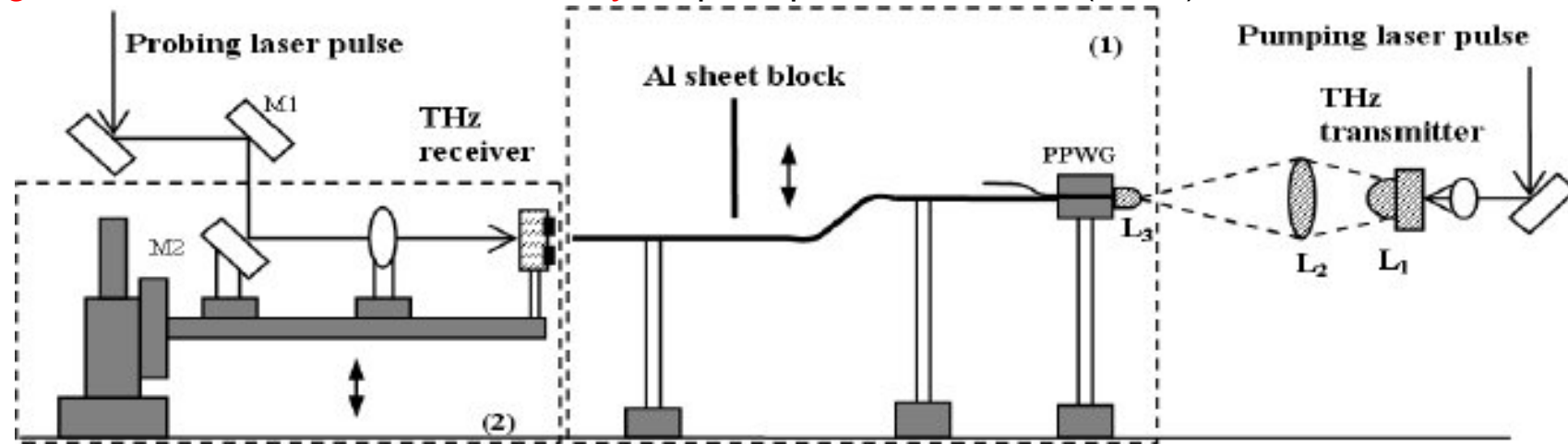


KAERI experiment layout



experimental observation of THz SEW—continue

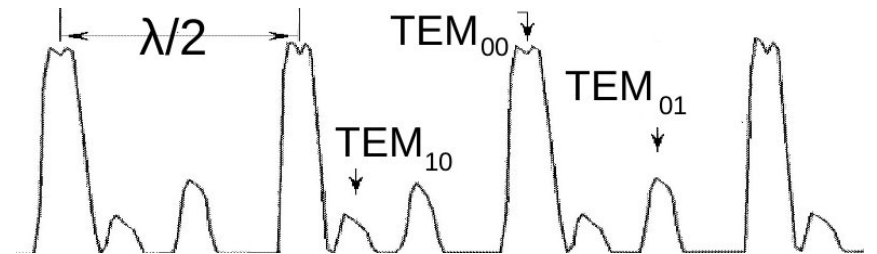
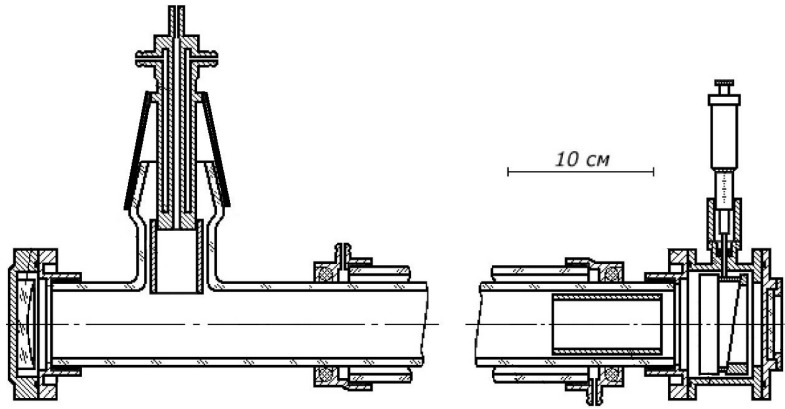
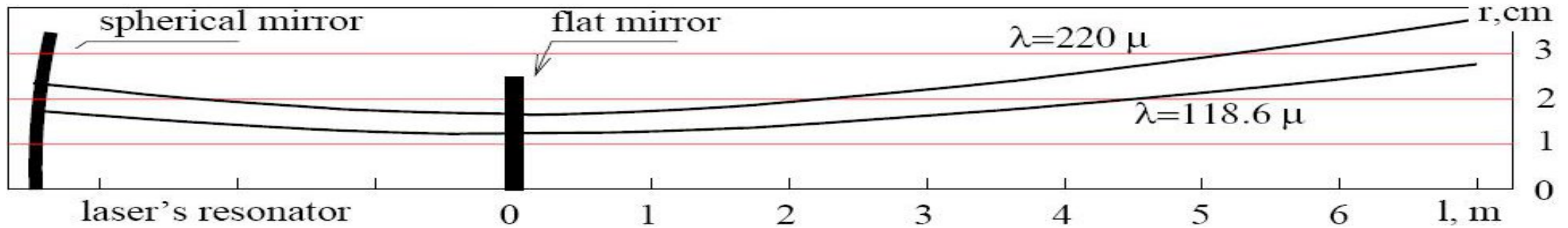
Mufei Gong¹, Tae-In Jeon², D. Grischkowsky, Opt.Express¹⁷, 17088 (2009)



H_2O laser

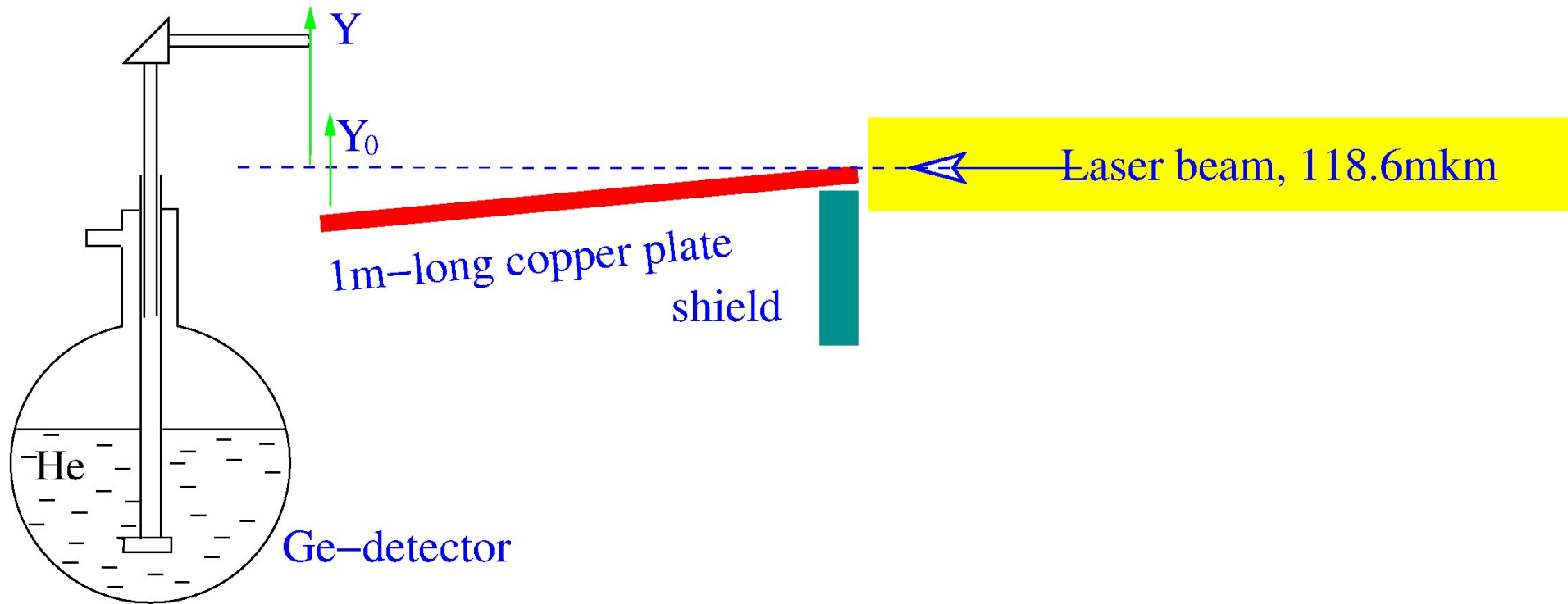
H_2O+H_2 : 220.2, 118.6, 79.1, 78.4, 55.1, 33, 28 μm ;

D_2O+D_2 : 171.7, 107.1, 84 μm

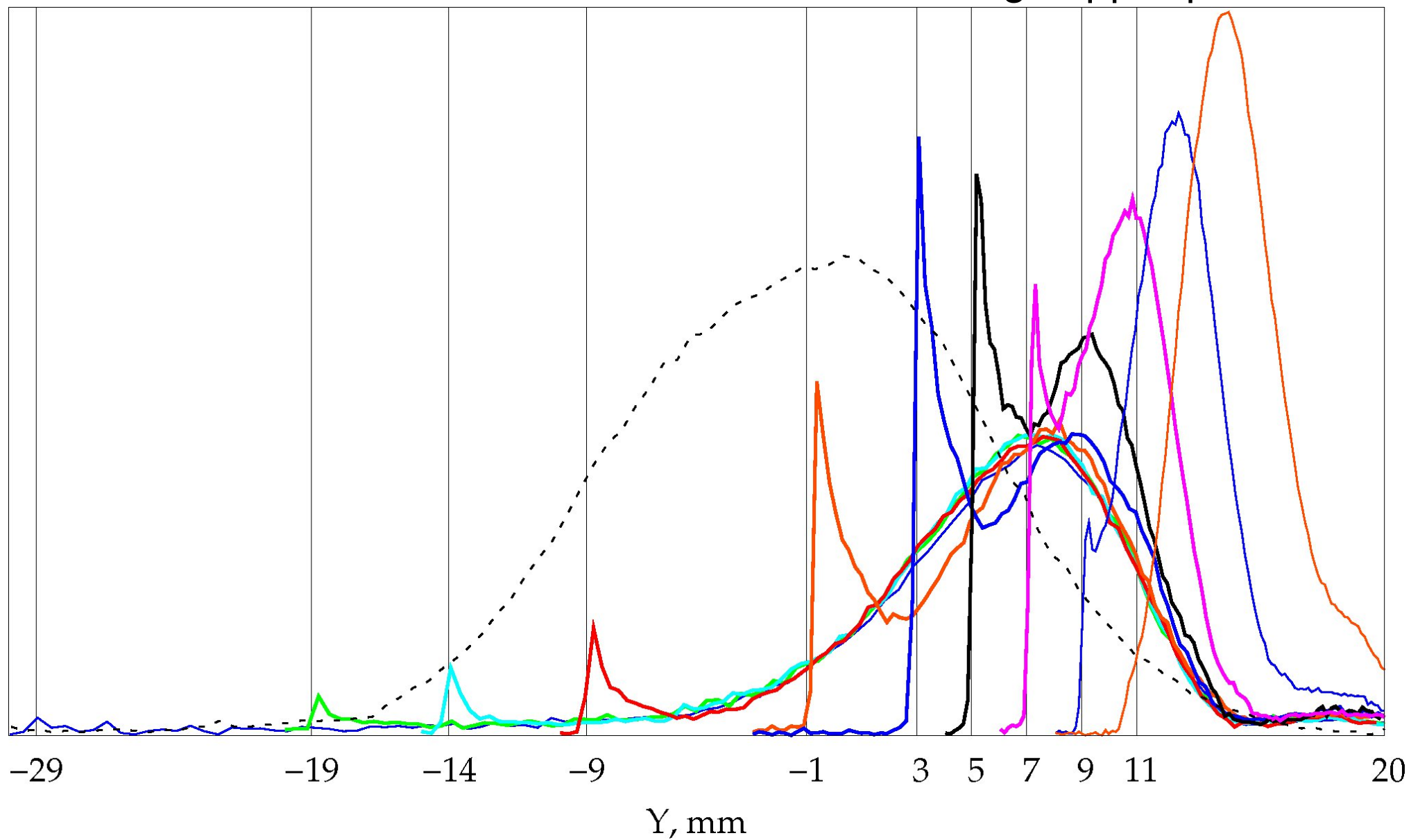


Our new experimental setup

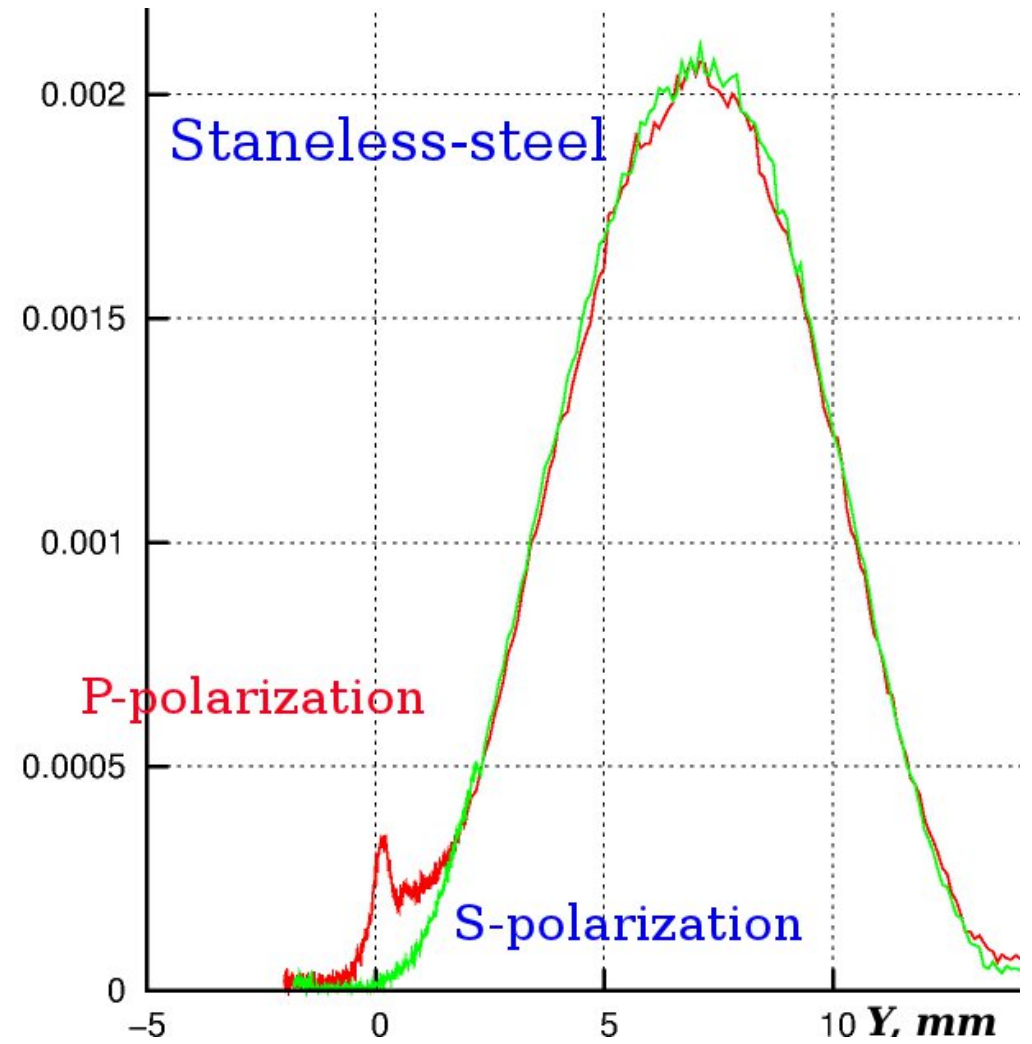
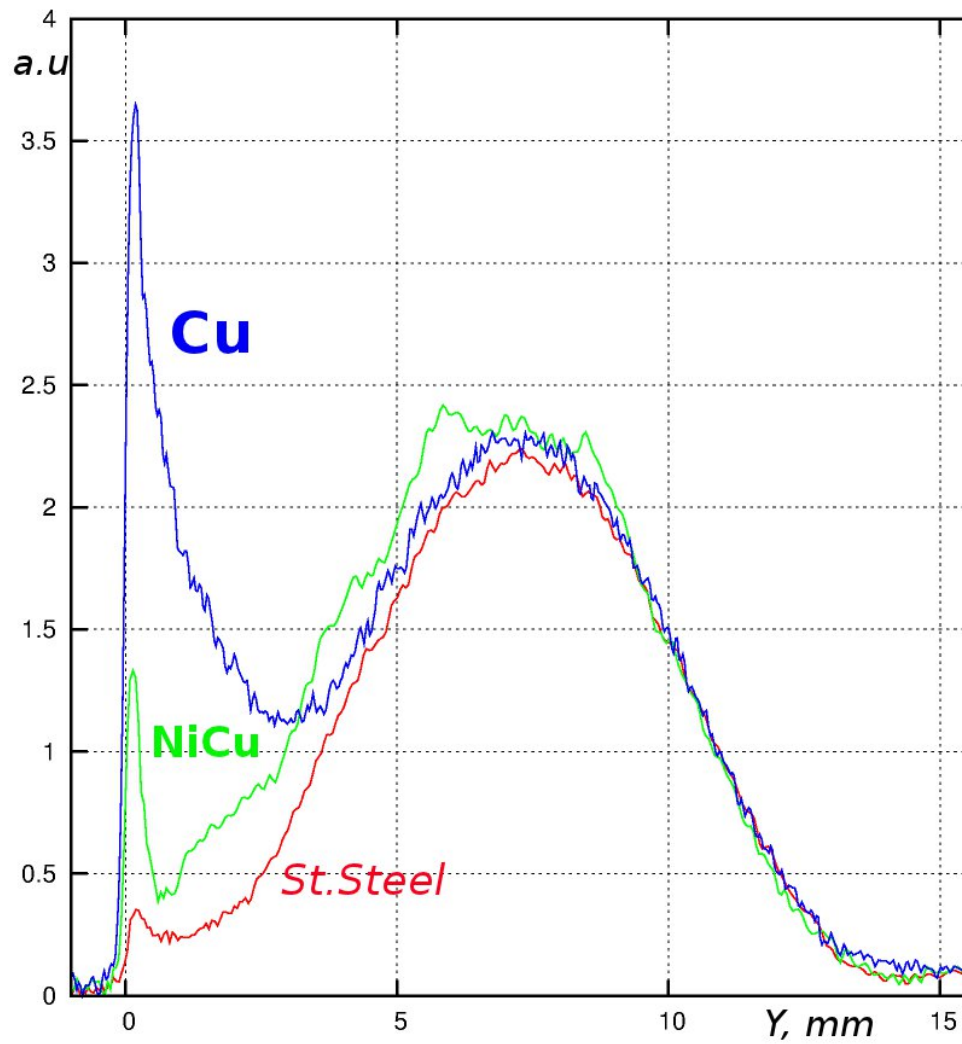
to measure Thz-beam cross-section at the end of the long copper plate

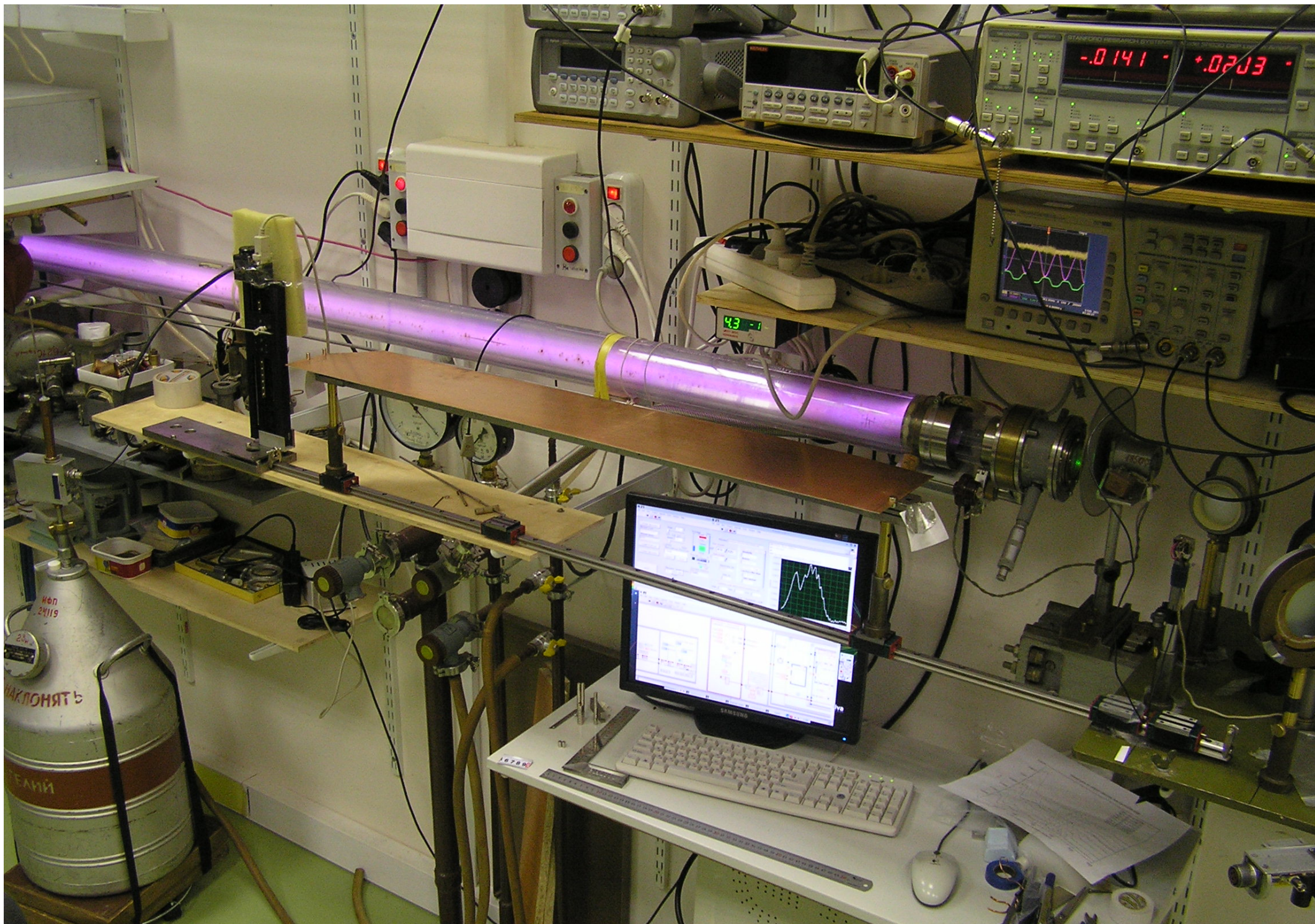
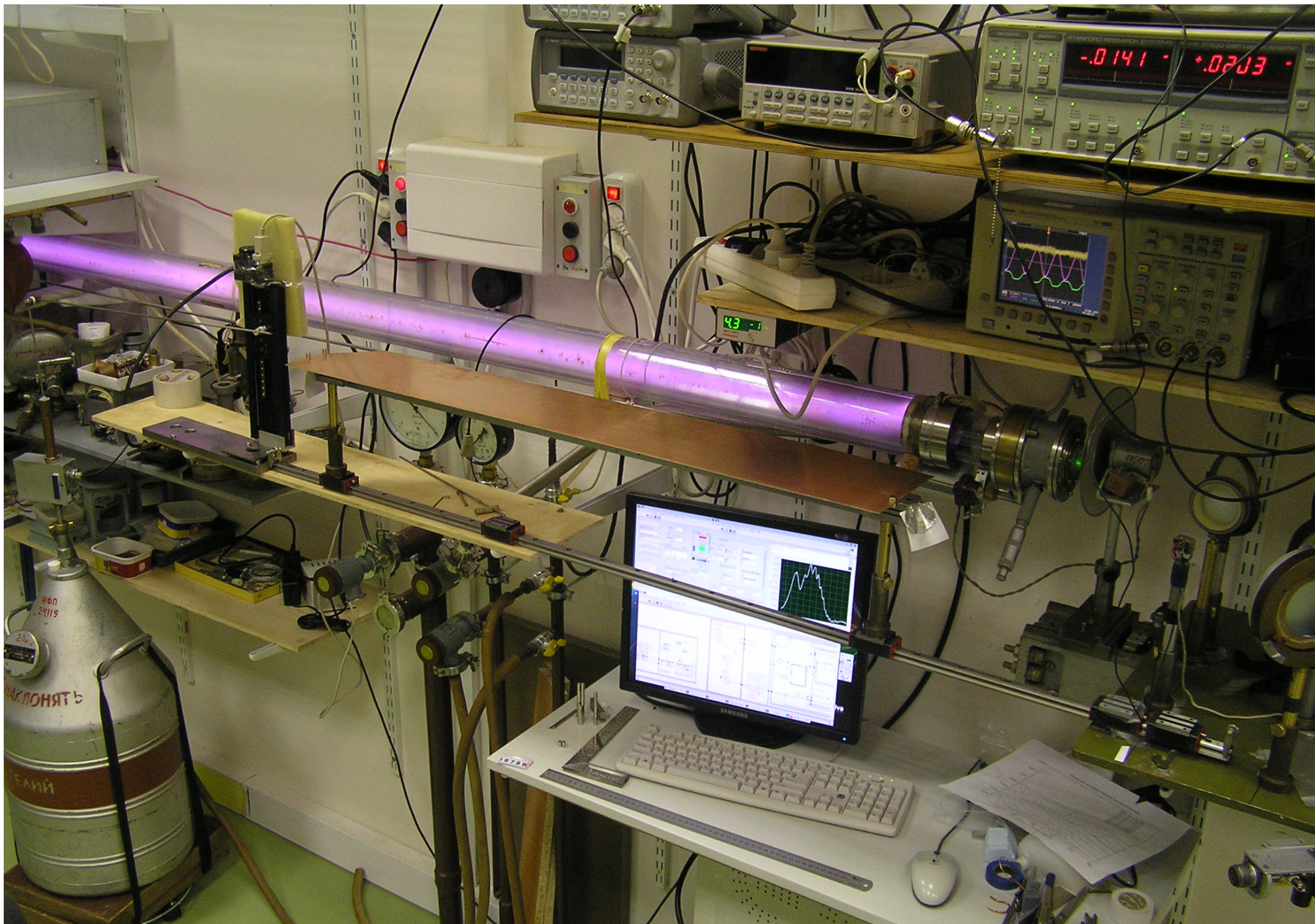


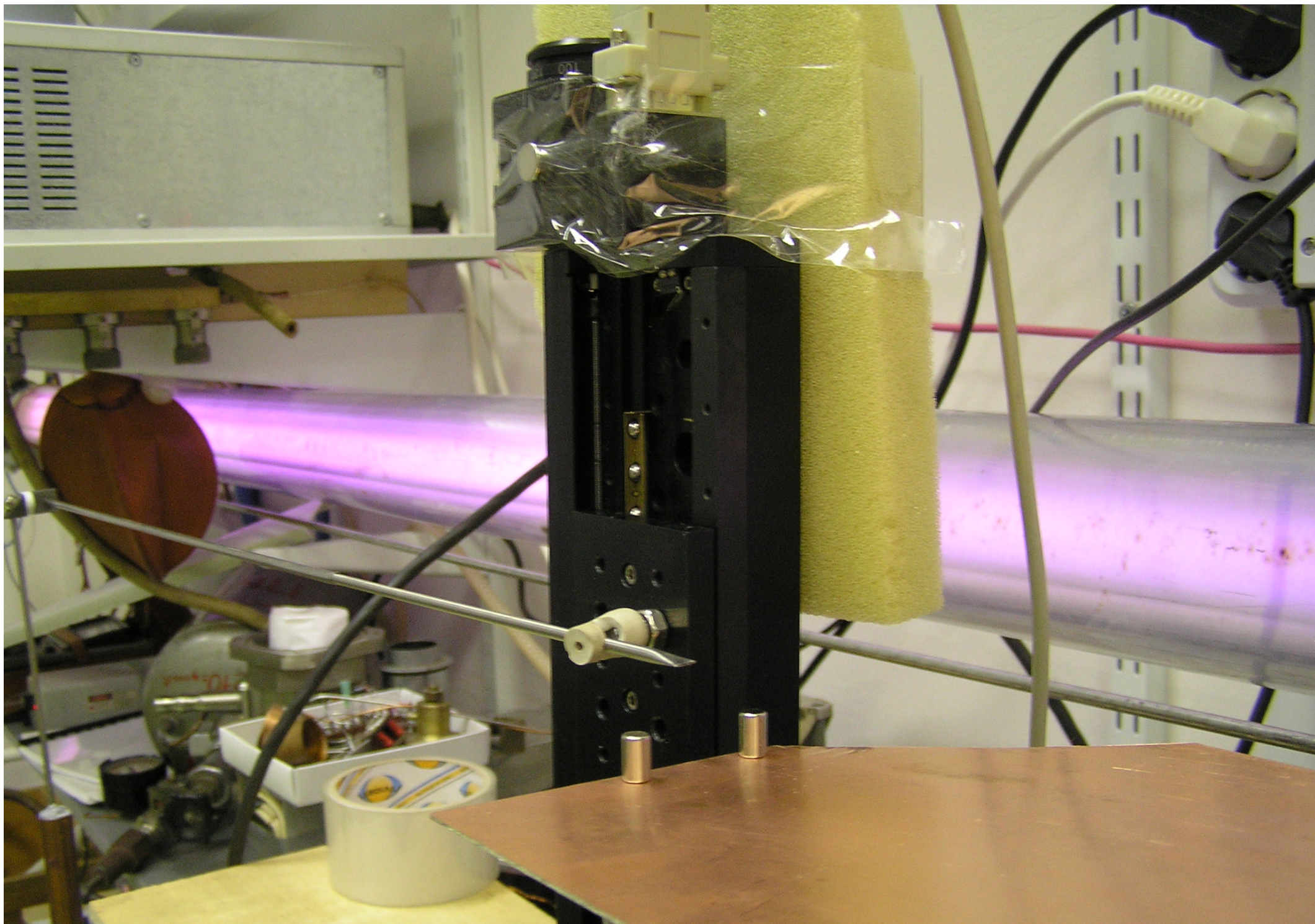
Thz-beam cross-section at the end of the long copper plate



Metal-dependant and polarization-dependant effects







RESULTS

- High-efficient excitation of THz-SEW on metal was achieved.
- The exponential damping of THz-SEW above bare metals was measured, and it agree well with the theory.
- Found, that under the THz-region the difference of longitudinal fase-velocities of SEW and free EM-wave are of no importance:
 - fase-matching couplers and decouplers are useless;
 - stand-alone SEW not possible;
 - to exite the SEW with the light beam, one has to match once their EM-field distributions.